

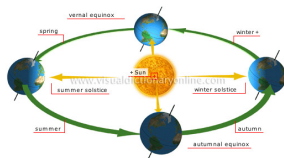
# Peculiar seasoning in the neutrino day-night asymmetry: where and when to look for spices?

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# What is seasoning?



## 1 Culinary seasoning

## 2 Solar $\nu$ seasonal effects & other time-regular effects

- $\sim 7\%$  flux variations due to seasonal variation of **Sun-to-Earth distance**
- solar cycles (including those assumed to exist due to **acoustic waves in the core**)  $\rightarrow$  neutrino flux variation
- Solar  $\nu$  **oscillations inside the Earth**  $\rightarrow$  **Day-Night Asymmetry** and its bizarre time variations (**OUR TALK!**)

## Why Solar $\nu$ Day-Night Asymmetry?

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- **Uniquely predictable** given the Solar and the Earth's structure and neutrino mixing params
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N.B. Extraction of DNA needs a long-term observation  $\rightarrow$  we inevitably face seasonal effects!

# Oscillations in matter & Earth regeneration [1]



The theory of DNA is quite conventional,

$$i\lambda\partial_x R(x, x_0) = H(x)R(x, x_0), \quad R(x_0, x_0) = 1;$$
$$H(x) = \left( -\cos 2\theta_0 + \frac{2EV(x)}{\Delta m^2} \right) \sigma_1 + \sin 2\theta_0 \sigma_3,$$

$R_{f,f'}(x, x_0) \equiv \langle \nu_f(x) | \nu_{f'}(x_0) \rangle$  is the flavor evolution matrix ( $f, f' = e, \mu$ )

$V(x) = \sqrt{2}G_F N_e(x)$  is the Wolfenstein potential  
 $\lambda = \Delta m^2 / 4E = \pi / \ell_{\text{osc}}, \quad \ell_{\text{osc}} \sim 20 \dots 300 \text{ km}$   
 $\sin^2 2\theta_0 \approx 0.86, \Delta m^2 \approx 7.6 \times 10^{-5} \text{ eV}^2$  [PDG2012]

$N_e(x)$  is the electron density  
 $E$  is the  $\nu$  energy  
 $x$  goes along the  $\nu$  ray



## Oscillations in matter & Earth regeneration [2]



$$i\lambda\partial_x R(x, x_0) = \left\{ \left( -\cos 2\theta_0 + \frac{2EV(x)}{\Delta m^2} \right) \sigma_1 + \sin 2\theta_0 \sigma_3 \right\} R(x, x_0)$$

- There are various **approximate** approaches to this equation which are relevant to the Earth regeneration effect for solar neutrinos [D'Olivo,1992; Supanitsky,D'Olivo,Medina-Tanco,2008; Lisi,Montanino,1997; de Holanda,Wei Liao,Smirnov,2004; Ioannian,Smirnov,2004; Blennow,Ohlsson,2004; Aleshin,Kharlanov,Lobanov,2013]
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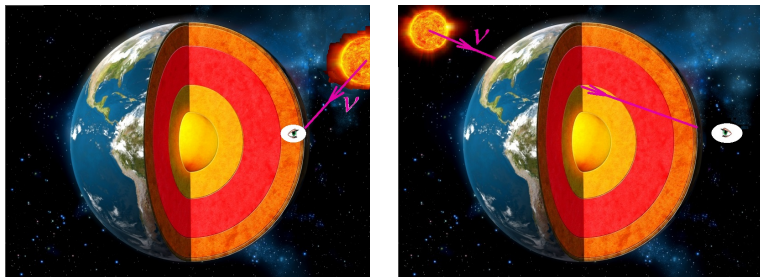


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These approaches give virtually the same results for solar neutrinos in the Earth, considered as a spherically-symmetric layered structure (**PREM**)

# Oscillations in matter & Earth regeneration [3]



Day/night probabilities of observing  $\nu_e$ :

$$P_e(\text{day}) = \frac{1}{2} + \frac{1}{2} \cos 2\theta_{\text{Sun}} \cos 2\theta_0,$$

$$P_e(\text{night}) = \frac{1}{2} + \frac{1}{2} \cos 2\theta_{\text{Sun}} \left\{ \cos 2\theta_n^- + 2 \sin 2\theta_0 \sum_{j=1}^{n-1} \Delta\theta_j \cos 2\Delta\psi_{n,j} \right\},$$

$n$  is the number of crossed interfaces between the Earth's layers

$\Delta\psi_{n,j} \approx \pi L_{n,j} / \ell_{\text{osc}}$  is the osc. phase diff. (detector- $j$ th crossing pt.)

$\Delta\theta_j$  are jumps of the effective mixing angle

$\theta_n^-$  is the effective mixing angle under the detector

As time goes by...

$$P_e(\text{night}) = \frac{1}{2} + \frac{1}{2} \cos 2\theta_{\text{sun}} \left\{ \cos 2\theta_n^- + 2 \sin 2\theta_0 \sum_{j=1}^{n-1} \Delta\theta_j \cos 2\Delta\psi_{n,j} \right\}$$

- The number of crossed interfaces  $n$  changes, depending on the **zenith angle**  $\Theta_Z(t)$
- The distance from the  $j$ th crossing pt. to the detector  
 $L_{n,j} = L_{n,j}(\Theta_Z(t))$

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In order to cope with such an object, we use the **stationary phase approximation**

$$\int_a^b F(\tau) e^{i\lambda S(\tau)} d\tau = \sqrt{\frac{2\pi i}{\lambda S''(\tau_0)}} F(\tau_0) e^{i\lambda S(\tau_0)} + \frac{F(\tau) e^{i\lambda S(\tau)}}{i\lambda S'(\tau)} \Big|_a^b + O(\lambda^{-3/2}), \quad \lambda \rightarrow +\infty,$$

where  $F(\tau)$  and  $S(\tau)$  are smooth on  $[a, b]$  and  $S'(\tau) = 0$  only at  $\tau = \tau_0 \in (a, b)$ .

It is easy to see that...

$$P_e(\text{night}) = \frac{1}{2} + \frac{1}{2} \cos 2\theta_{\text{sun}} \left\{ \cos 2\theta_n^- + 2 \sin 2\theta_0 \sum_{j=1}^{n-1} \Delta\theta_j \cos 2\Delta\psi_{n,j} \right\}$$

- Stationary points  $\Theta'_Z(t) = 0$ : **midnights** when integrating over the night and **solstices** when integrating over the seasons;
- Edge terms vanish
- We take the dependence  $\Theta_Z(t)$  from **spherical astronomy**
- The **small parameter** here is  $\ell_{\text{osc}}/L_{n,j}$



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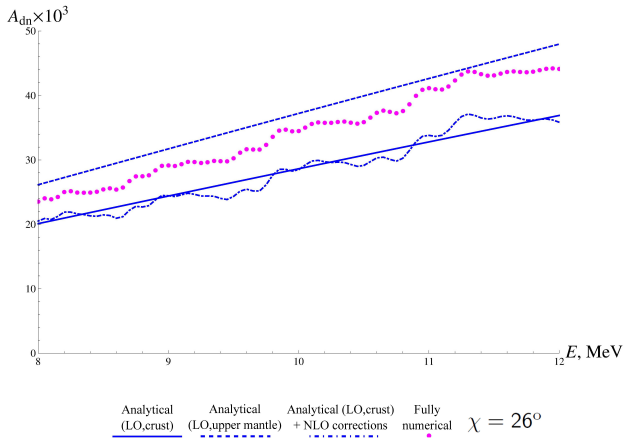
$$\langle P_e(\text{night}) - P_e(\text{day}) \rangle_{\text{year}} \approx \frac{1}{2} \cos 2\theta_{\text{Sun}} (\cos 2\theta_n^- - \cos 2\theta_0)$$

$$+ \cos 2\theta_{\text{Sun}} \sin 2\theta_0 \sum_{j=1}^{n-1} \Delta\theta_j \sum_{s=\pm 1} \frac{\vartheta(r_j - r_n \sin(\chi + s\varepsilon))}{2\pi \sqrt{\sin \varepsilon \cos \chi \sin(\chi + s\varepsilon)}} \frac{\sqrt{r_j^2/r_n^2 - \sin^2(\chi + s\varepsilon)}}{\lambda L_{n,j}^{\text{solstice}}} \times \cos\{2\Delta\psi_{n,j}^{\text{solstice}} + s'(s-1)\pi/4\},$$

$s' \equiv \text{sgn}\{L_{n,j}^{\text{solstice}} - r_n \cos(\chi + s\varepsilon)\}$ ,  $\Theta_Z = \pi - \chi - s\varepsilon$  (**solstice**  $s = \pm 1$ );  $\chi$  is the detector latitude,  $\varepsilon = 23.5^\circ$ ;  $r_j$  are the radii of the Earth's shells.

# A test drive: analytics vs. numerics

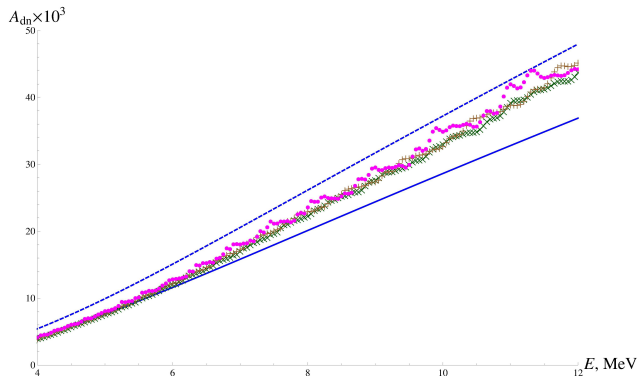
$$A_{dn} = \frac{2(P_e(\text{night}) - P_e(\text{day}))}{P_e(\text{night}) + P_e(\text{day})}$$



The constant vertical shift is due to the unaccounted fine structure of the crust under the detector. It is smooth enough and is season-independent, so we do not bother about it.

# A test drive: different latitudes

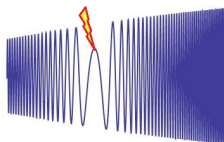
$$A_{\text{dn}} = \frac{2(P_e(\text{night}) - P_e(\text{day}))}{P_e(\text{night}) + P_e(\text{day})}$$



Gran Sasso Kamioka Tropic  
(42.5N) (36.4N) (23.5N)  
x + •

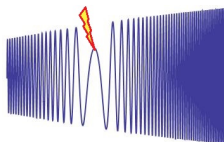
Indeed, the closer to the Tropic, the more vivid are the oscillatory contributions of the stationary points!

# The Miracles of the Stationary points



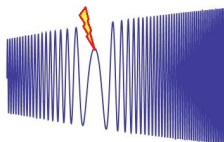
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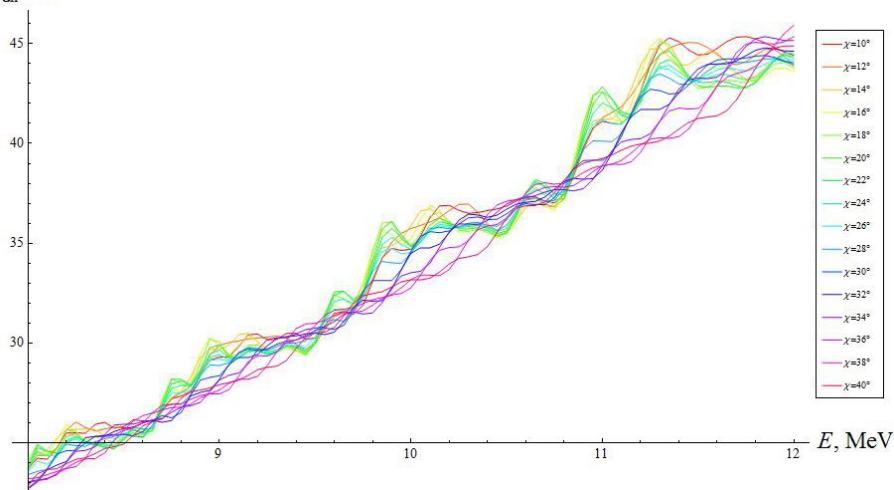


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- Moreover, it is wrong to say that *the Earth's core does not contribute to DNA since the Sun rarely descends low enough to shine through it.* Rareness is not a measure for this localized contribution!

# The tropical Sun

$$A_{\text{dn}} = \frac{2(P_e(\text{night}) - P_e(\text{day}))}{P_e(\text{night}) + P_e(\text{day})}$$

$A_{\text{dn}} \times 10^3$



The curves around  $\chi = 23.5^\circ$  exhibit the clean and vivid oscillations; the positions of the peaks are **very** sensitive to  $\Delta m^2$  and  $\{r_i\}$

# Interference experiment in $\nu$ oscillations?



- **When?** Near the **solstices**
- **Where?** Near the Tropics, the closer the better (Sao Paolo, **23°33'S**)

# When and where to look for spices? (a conclusion)

- **When?** Near the **solstices**
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- **Why?** To extremely precisely determine the **radii of the Earth's shells** and the solar neutrino **mass-squared difference**

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- **When?** Near the **solstices**
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- **Why?** To extremely precisely determine the **radii of the Earth's shells** and the solar neutrino **mass-squared difference**
- **Really?** In order to distinguish these effects, one needs **within one order more detection events** at the currently achieved energy resolution  $\delta E \sim 0.5$  MeV, if one employ the **adaptive recognition** of wave-like patterns on the  $A_{dn}(E)$  profile [**to be published**]

## Acknowledgments & publications

The numerical simulations were made using the Supercomputing cluster “Lomonosov” (MSU)

### References

- [1] S. S. Aleshin, O. G. Kharlanov, and A. E. Lobanov, Analytical treatment of long-term observations of the day-night asymmetry for solar neutrinos, *Phys. Rev. D* **87**, 045025 (2013).
- [2] O. G. Kharlanov, and A. E. Lobanov, Peculiar seasonal effects in the neutrino day-night asymmetry, submitted to *Phys. Rev. D*.

Thank you for your attention!