

# Production of kink–antikink pair in collisions of high–energy particles

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# Topological solitons in weakly coupled theories

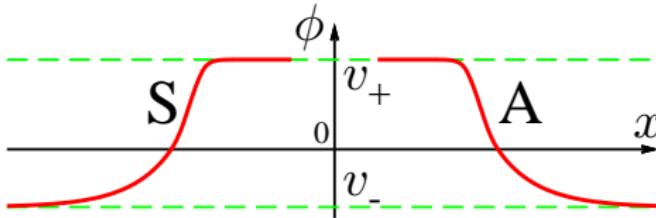
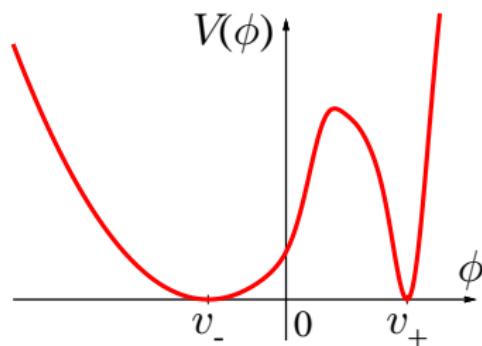
$\boxed{\phi}$

(1 + 1)

$\hbar = c = 1$

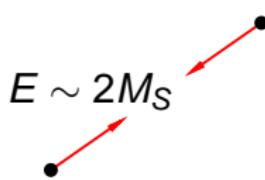
$$S = \frac{1}{g^2} \int dx dt \left[ (\partial_\mu \phi)^2 / 2 - V(\phi) \right]$$

- semiclassical parameter
- coupling constant ( $\phi = g\phi'$ )



Properties:  $L_S \sim m^{-1}$   
 $M_S \sim m/g^2$

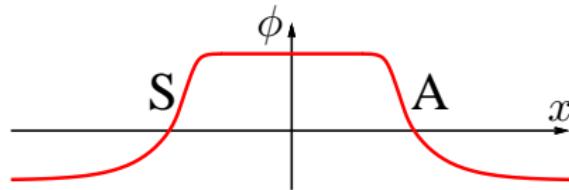
# From particles to solitons



$$E \sim 2M_S$$

$$\lambda \sim 1/E \sim g^2/m$$

Exponential suppression!



$$L_S \sim 1/m$$

$$\mathcal{P}(E) \approx A \cdot e^{-F(E)/g^2}$$

- Coherent-state “estimate”

$$\bar{n}_S \sim M_S/m \sim 1/g^2$$

$$\langle 2|SA \rangle \sim e^{-\bar{n}_S/2} \sim e^{-c/g^2}$$

Drukier, Nussinov (1982)

Banks et al. (1990)

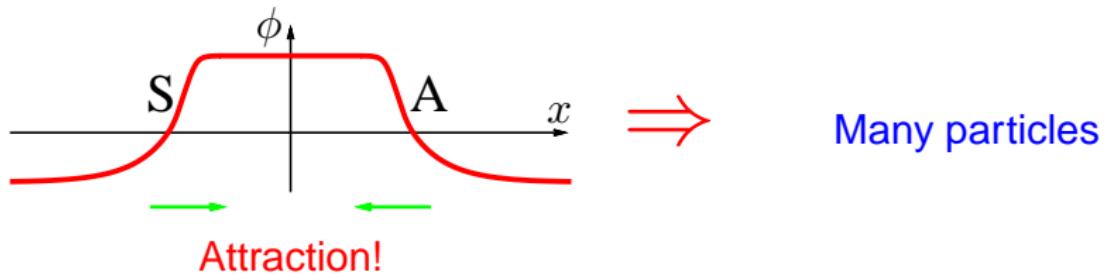
- Unitarity arguments

Zakharov (1991)

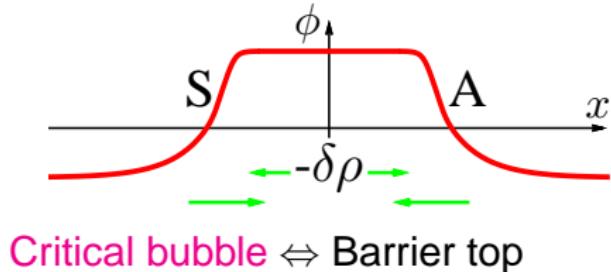
No reliable estimate of  $\mathcal{P}(E)$  so far!

Aim: calculate semiclassically  $F(E)$ .

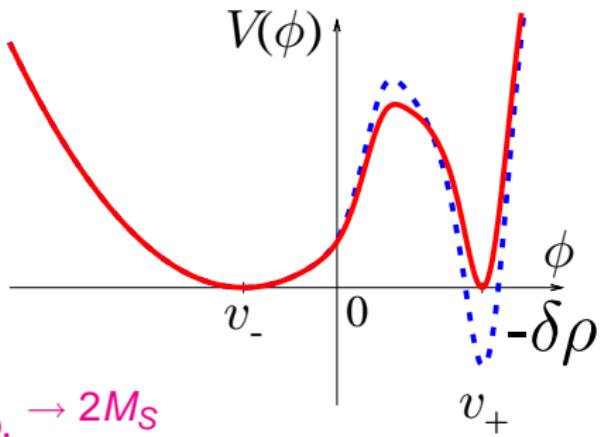
# Introducing potential barrier



Not a tunneling process?



$\delta\rho \rightarrow 0$ : cr. bubble  $\rightarrow$  SA;  $E_{\text{cr.b.}} \rightarrow 2M_S$



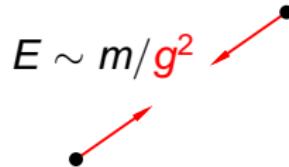
## In-state

Rubakov, Son, Tinyakov, 1992

Not semiclassical!

RST conjecture:  $F(E)$  universal

Does not depend on details of the in-state



$$\mathcal{P}(E, N) = \sum_{i,f} \left| \langle i | \hat{P}_E \hat{P}_N \hat{S} | f \rangle \right|^2$$

↑  
Projectors

- $N \gg 1 \Rightarrow$  semiclassical in-states
- $N \ll 1/g^2 \Rightarrow F(E, N) \approx F(E)$

Checks of universality:

- Field theory

Tinyakov, 1991

Mueller, 1992

- Toy QM models

Bonini et al, 1999

Levkov et al, 2009

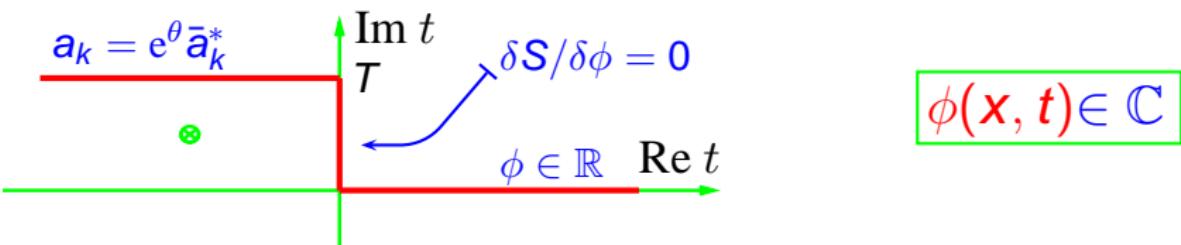
$$F(E) = \lim_{g^2 N \rightarrow 0} F(E, N)$$

# Semiclassical method

Rubakov, Son, Tinyakov, 1992

$$\mathcal{P}(E, N) = \sum_{i,f} \left| \langle i | \hat{P}_E \hat{P}_N \hat{S} | f \rangle \right|^2 = \int [d\phi][d\phi'] e^{iW}$$

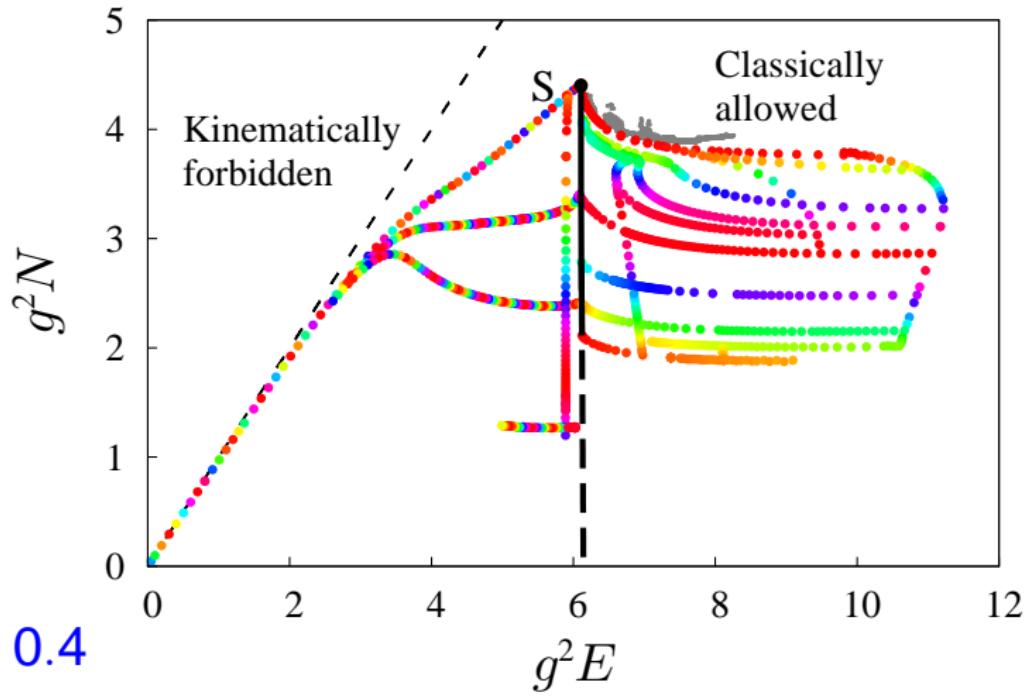
$W \propto 1/g^2 \Rightarrow$  Saddle-point method!



$$\mathcal{P} = A \cdot e^{-F/g^2} \quad \frac{1}{g^2} F(E, N) = 2\text{Im}S[\phi] - TE - \theta N$$

# Numerical solutions

$$V(\phi) = \frac{1}{2}(\phi + 1)^2 \left[ 1 - v \cdot f\left(\frac{\phi-1}{a}\right) \right], \quad f(x) = e^{-x^2} (1 + x^3 + x^5)$$

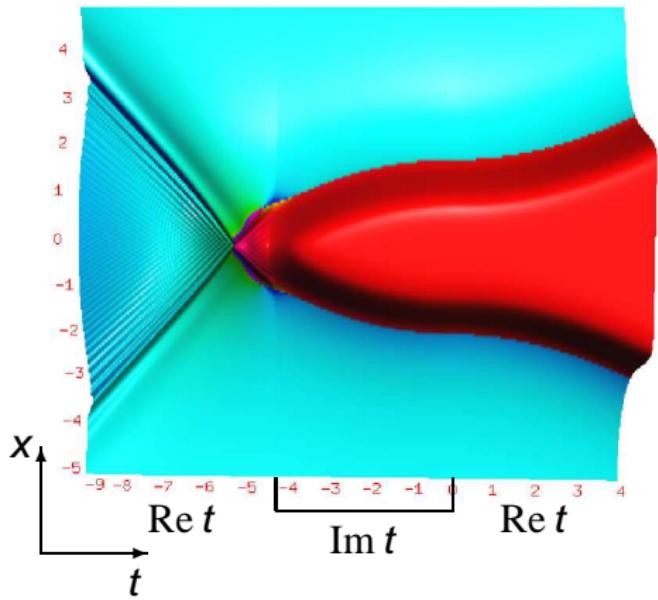


$$\delta\rho = 0.4$$

# Going to $E > 2M_S$

$$E \approx 5.48$$

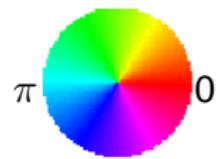
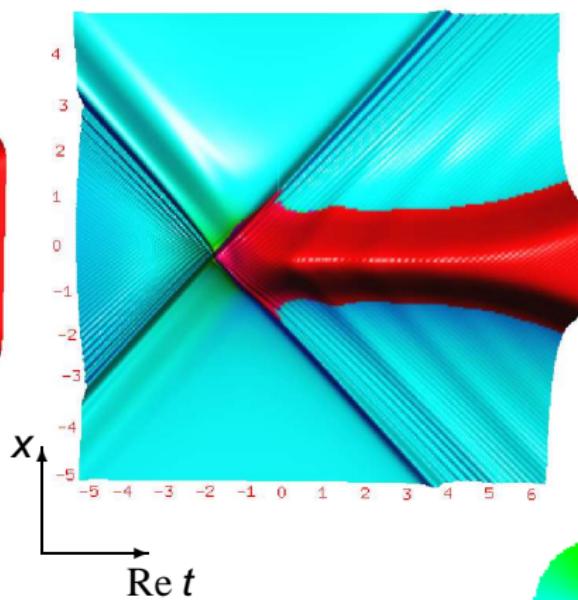
$$N \approx 2.39, \delta\rho = 0.4$$



$$(2M_S \approx 6.23)$$

$$E \approx 9.06$$

$$N \approx 2.47, \delta\rho = 0.4$$



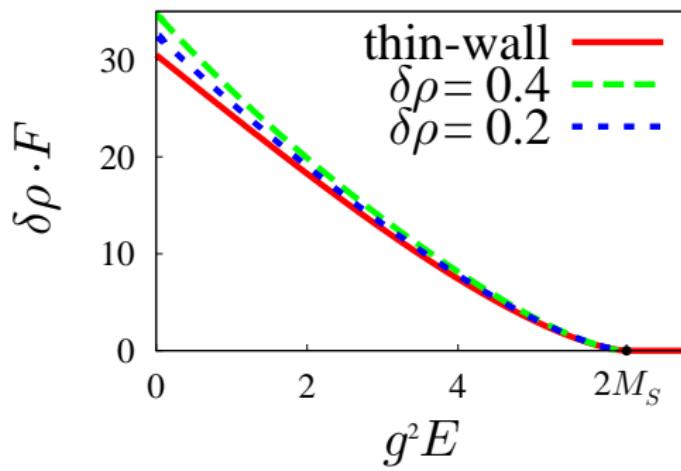
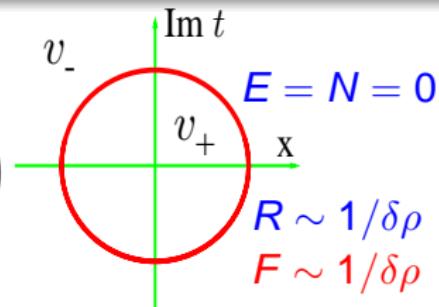
$\delta\rho \rightarrow 0$ : thin-wall limit!

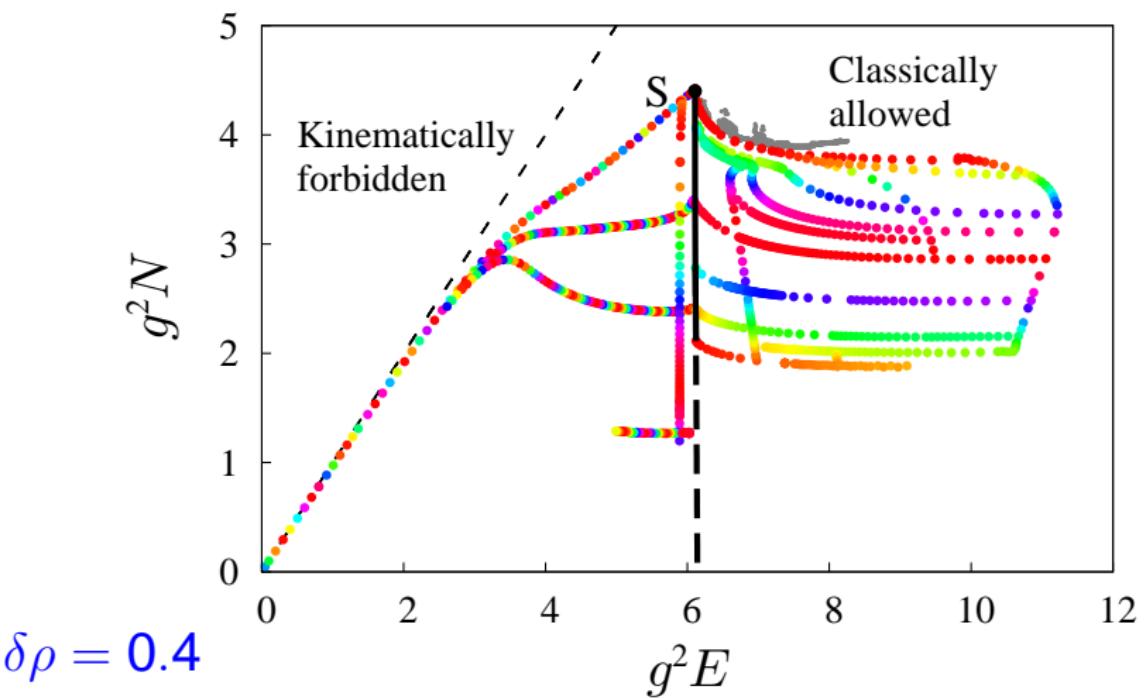
$$F(\delta\rho) = F_{-1}/\delta\rho + F_0 + O(\delta\rho)$$

$$F_{-1}(E, N) = E_S^2 \left( \pi - 2 \arcsin \frac{E}{2E_S} - \frac{E}{E_S} \sqrt{1 - \frac{E^2}{4E_S^2}} \right)$$

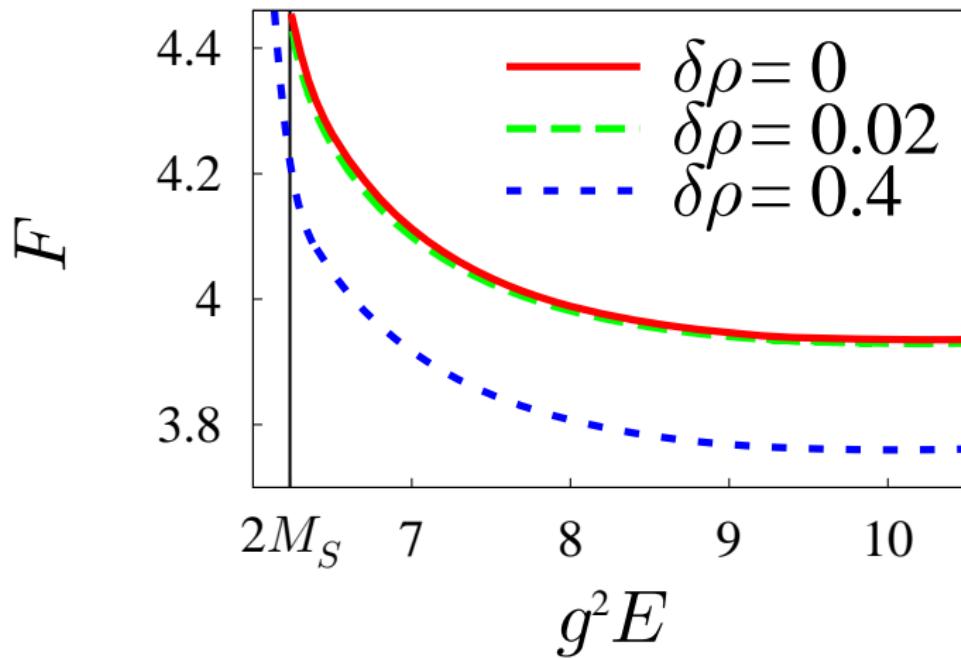
Voloshin, Selivanov, 1986

Rubakov et al, 1991



Extrapolating to  $N \rightarrow 0$ 

## Result



# Conclusions

- Method is applicable in  **$2D$**  scalar field models.
- The probability of **SA** creation in high-energy collisions is

$$\mathcal{P}(E) \approx e^{-F(E)/g^2}$$