

The muon g-2 discrepancy: errors or new physics?

Massimo Passera
INFN Padova

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The present experimental values:

$$a_e = 1159652180.73 (28) \times 10^{-12}$$

0.24 parts per billion !! Hanneke et al., PRL100 (2008) 120801

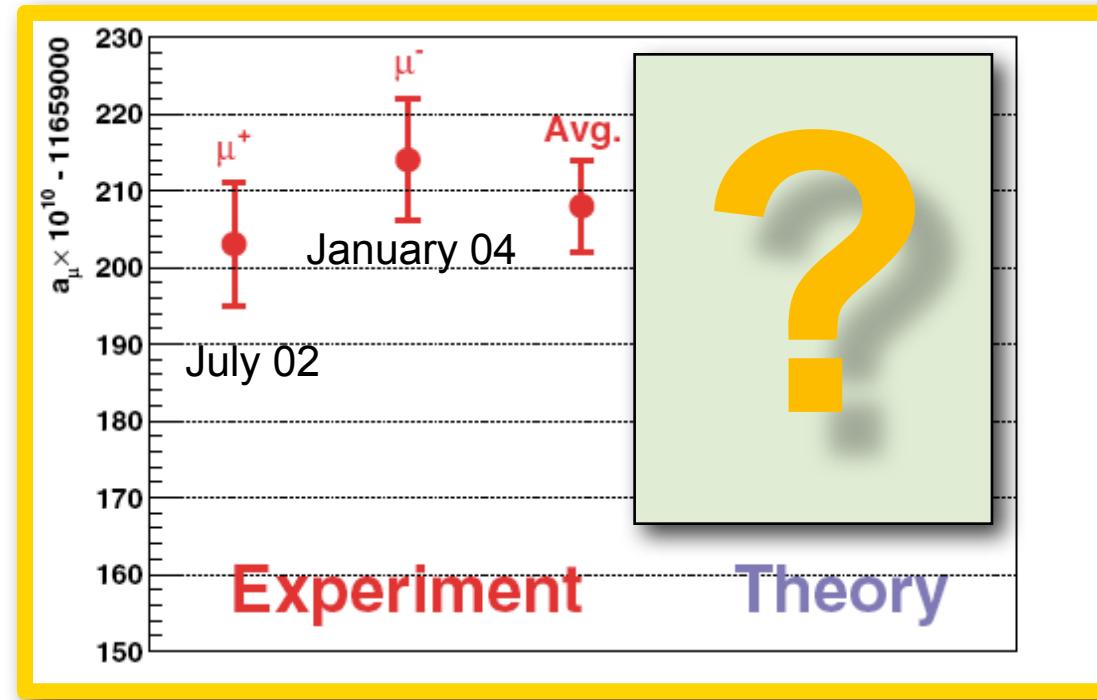
$$a_\mu = 116592080 (63) \times 10^{-11}$$

0.5 parts per million !! E821 - Final Report: PRD73 (2006) 072003

$$a_\tau = -0.018 (17)$$

DELPHI - EPJC35 (2004) 159 [$a_\tau^{\text{SM}} = 117721(5) \times 10^{-8}$, Eidelman & MP '07]

The muon $g-2$: experimental result



- Today: $a_\mu^{\text{EXP}} = (116592080 \pm 54_{\text{stat}} \pm 33_{\text{sys}}) \times 10^{-11}$ [0.5 ppm].
- Future: a new $(g-2)_\mu$ exp aims at 0.14 ppm! [D. Hertzog, KLOE2 Physics Workshop, LNF, April 2009].
- Are theorists ready for this? [not yet]

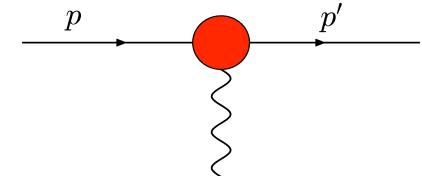
- The Dirac theory predicts for a lepton $l=e,\mu,\tau$:

$$\vec{\mu}_l = g_l \left(\frac{e}{2m_l c} \right) \vec{s} \quad g_l = 2$$

- QFT predicts deviations from the Dirac value:

$$g_l = 2(1 + a_l)$$

- Study the photon-lepton vertex:



$$\bar{u}(p') \Gamma_\mu u(p) = \bar{u}(p') \left[\gamma_\mu F_1(q^2) + \frac{i\sigma_{\mu\nu}q^\nu}{2m} F_2(q^2) + \dots \right] u(p)$$

$$F_1(0) = 1 \quad F_2(0) = a_l$$

The QED contribution to a_μ

$$a_\mu^{\text{QED}} = (1/2)(\alpha/\pi) \quad \text{Schwinger 1948}$$

$$+ 0.765857408 (27) (\alpha/\pi)^2$$

Sommerfield; Petermann; Suura & Wichmann '57; Elend '66; MP '04

$$+ 24.05050959 (42) (\alpha/\pi)^3$$

Remiddi, Laporta, Barbieri ... ; Czarnecki, Skrzypek; MP '04;

Friot, Greynat & de Rafael '05, Mohr, Taylor & Newell '08

$$+ 130.805 (8) (\alpha/\pi)^4$$

Kinoshita & Lindquist '81, ... , Kinoshita & Nio '04, '05;

Aoyama, Hayakawa, Kinoshita & Nio, June & Dec 2007

$$+ 663 (20) (\alpha/\pi)^5 \quad \text{In progress}$$

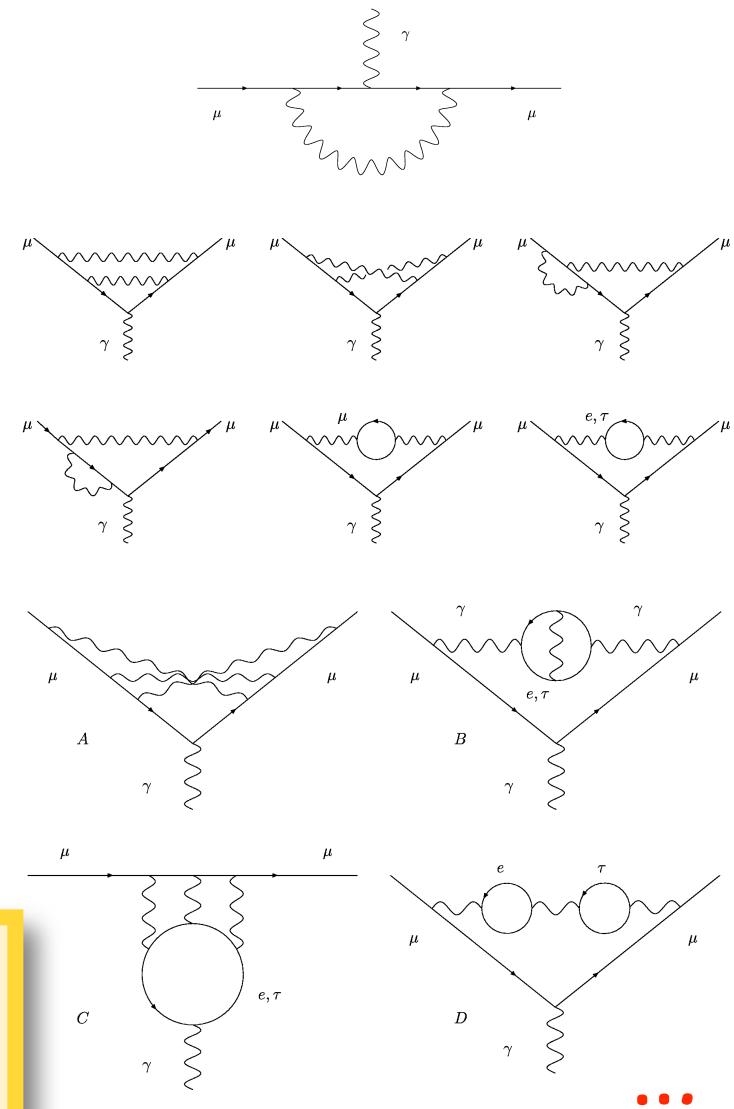
Kinoshita et al. '90, Yelkhovsky, Milstein, Starshenko, Laporta, Karshenboim, ..., Kataev, Kinoshita & Nio March '06.

Adding up, I get:

$$a_\mu^{\text{QED}} = 116584718.08 (14)(04) \times 10^{-11}$$

from coeffs, mainly from 5-loop unc from new $\delta\alpha$ ('08)

with $\alpha=1/137.035999084(51)$ [0.37 ppb]



[A parenthesis on the electron g-2...]

$$a_e^{\text{SM}}$$

$$= (1/2)(\alpha/\pi) - 0.328\,478\,444\,002\,89(60) (\alpha/\pi)^2$$

Schwinger 1948

Sommerfield; Petermann; Suura & Wichmann '57; Elend '66; MP '06

$$A_2^{(4)} (m_e/m_\mu) = 5.197\,386\,78(26) \times 10^{-7}$$

$$A_2^{(4)} (m_e/m_\tau) = 1.837\,62(60) \times 10^{-9}$$

$$+ 1.181\,234\,016\,827(19) (\alpha/\pi)^3$$

Kinoshita, Barbieri, Laporta, Remiddi, ..., Li, Samuel, Mohr, Taylor & Newell '08, MP '06

$$A_2^{(6)} (m_e/m_\mu) = -7.373\,941\,73(27) \times 10^{-6}$$

$$A_2^{(6)} (m_e/m_\tau) = -6.5819(19) \times 10^{-8}$$

$$A_3^{(6)} (m_e/m_\mu, m_e/m_\tau) = 1.909\,45(62) \times 10^{-13}$$

$$- 1.9144(35) (\alpha/\pi)^4$$

Kinoshita & Lindquist '81, ..., Kinoshita & Nio '05; Aoyama, Hayakawa, Kinoshita & Nio, June '07

$$+ 0.0(4.6) (\alpha/\pi)^5$$

In progress (12672 mass ind. diagrams!)

Aoyama, Hayakawa, Kinoshita, Nio '07; Aoyama, Hayakawa, Kinoshita, Nio & Watanabe 06/2008.

$$+ 1.676(20) \times 10^{-12} \text{ Hadronic}$$

Krause 1997, Jegerlehner & Nyffeler 2009

$$+ 0.02973(52) \times 10^{-12} \text{ Electroweak}$$

Mohr, Taylor & Newell '08; Czarnecki, Krause, Marciano '96

... and the best determination of alpha]

- The new measurement of the electron g-2 is:

$$a_e^{\text{exp}} = 1159652180.73 (28) \times 10^{-12} \text{ Hanneke et al, PRL100 (2008) 120801}$$

vs. old (factor of 15 improvement, 1.8σ difference):

$$a_e^{\text{exp}} = 1159652188.3 (4.2) \times 10^{-12} \text{ Van Dyck et al, PRL59 (1987) 26}$$

- Equating $a_e^{\text{SM}}(\alpha) = a_e^{\text{exp}}$ → best determination of alpha to date:

$$\alpha^{-1} = 137.035\ 999\ 084 (12)(37)(2)(33) [0.37\text{ppb}] \text{ Hanneke et al, '08}$$

δC_4^{qed} δC_5^{qed} δa_e^{had} $\delta a_e^{\text{exp}} \text{ (smaller than th!)}$

- Compare it with other determinations (independent of a_e):

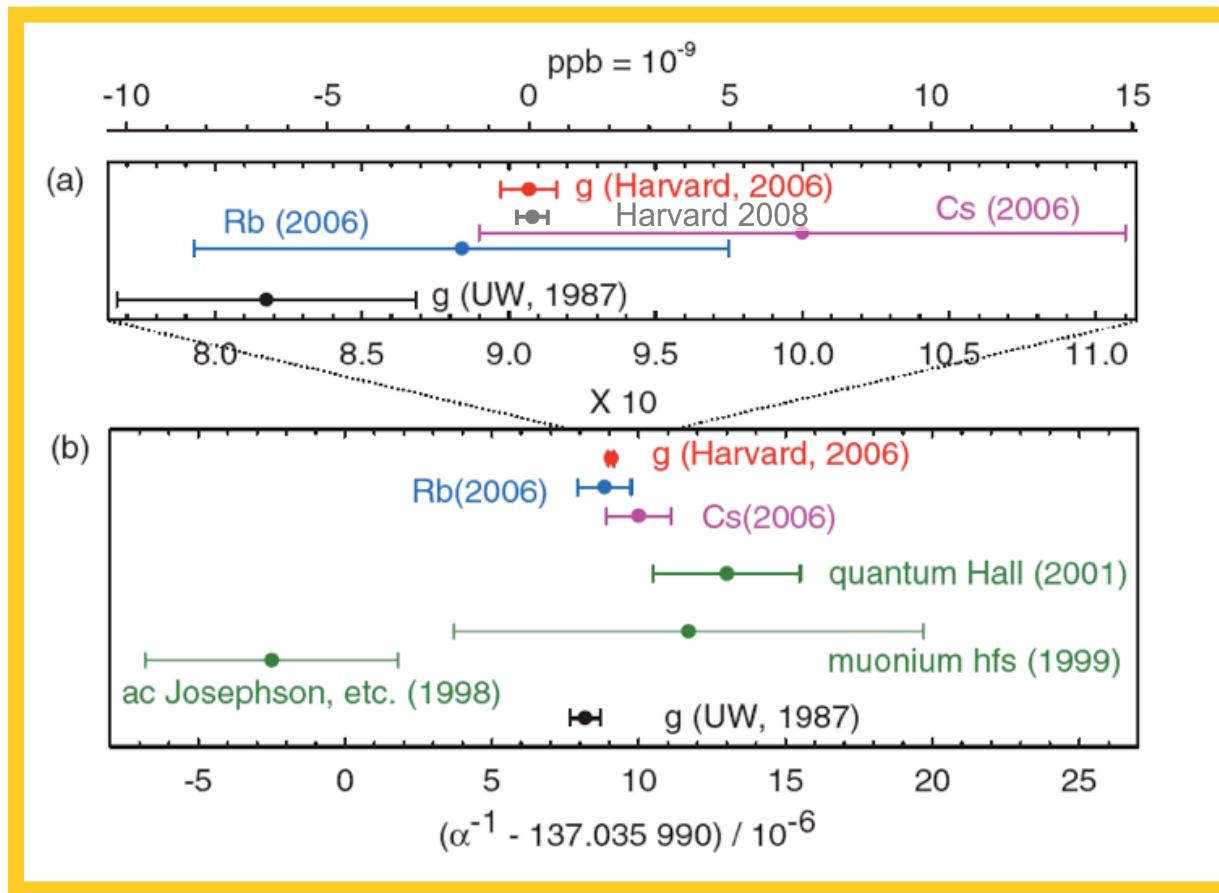
$$\alpha^{-1} = 137.036\ 000\ 00 \quad (110) \quad [7.7\text{ ppb}] \quad \text{PRA73 (2006) 032504 (Cs)}$$

$$\alpha^{-1} = 137.035\ 998\ 78 \quad (91) \quad [6.7\text{ ppb}] \quad \text{PRL96 (2006) 033001 (Rb)}$$

$$\alpha^{-1} = 137.035\ 999\ 45 \quad (62) \quad [4.6\text{ ppb}] \quad \text{PRL101 (2008) 230801 (Rb)}$$

Excellent agreement → beautiful test of QED at 4-loop level!

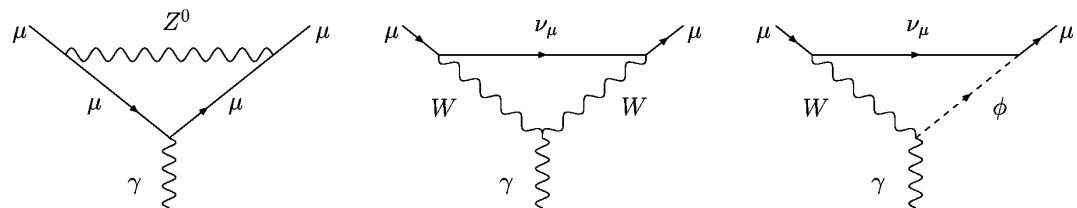
Old and new determinations of alpha



Gabrielse, Hanneke, Kinoshita, Nio & Odom, PRL99 (2007) 039902
Hanneke, Fogwell & Gabrielse, PRL100 (2008) 120801

The Electroweak contribution

One-loop term:



$$a_\mu^{\text{EW}}(\text{1-loop}) = \frac{5G_\mu m_\mu^2}{24\sqrt{2}\pi^2} \left[1 + \frac{1}{5} (1 - 4 \sin^2 \theta_W)^2 + O\left(\frac{m_\mu^2}{M_{Z,W,H}^2}\right) \right] \approx 195 \times 10^{-11}$$

1972: Jackiw, Weinberg; Bars, Yoshimura; Altarelli, Cabibbo, Maiani; Bardeen, Gastmans, Lautrup; Fujikawa, Lee, Sanda; Studenikin et al. '80s

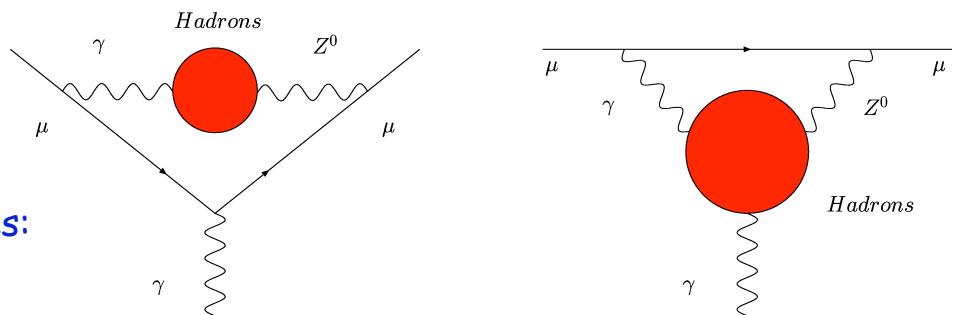
One-loop plus higher-order terms:

$$a_\mu^{\text{EW}} = 154 (2) (1) \times 10^{-11}$$

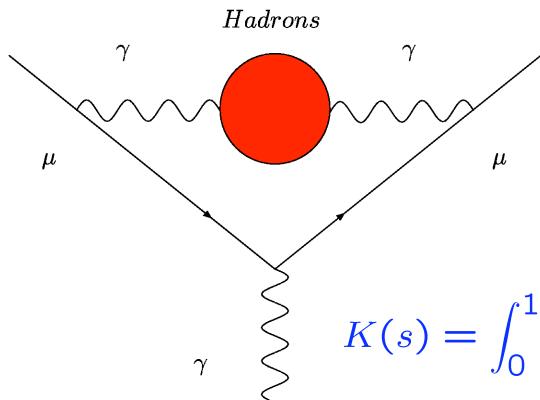
Higgs mass, M_{top} error,
3-loop nonleading logs

Hadronic loop uncertainties:

Kukhto et al. '92; Czarnecki, Krause, Marciano '95; Knecht, Peris, Perrottet, de Rafael '02; Czarnecki, Marciano, Vainshtein '02; Degrassi, Giudice '98; Heinemeyer, Stockinger, Weiglein '04; Gribouk, Czarnecki '05; Vainshtein '03.



The hadronic leading-order (HLO) contribution

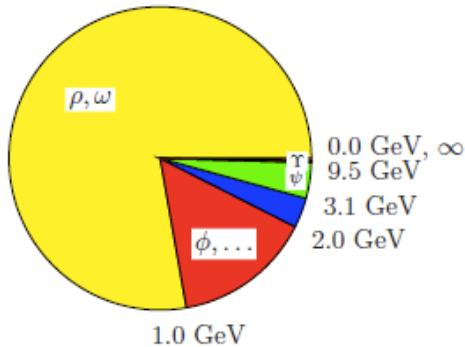


$$a_\mu^{\text{HLO}} = \frac{1}{4\pi^3} \int_{4m_\pi^2}^\infty ds K(s) \sigma^{(0)}(s) = \frac{\alpha^2}{3\pi^2} \int_{4m_\pi^2}^\infty \frac{ds}{s} K(s) R(s)$$

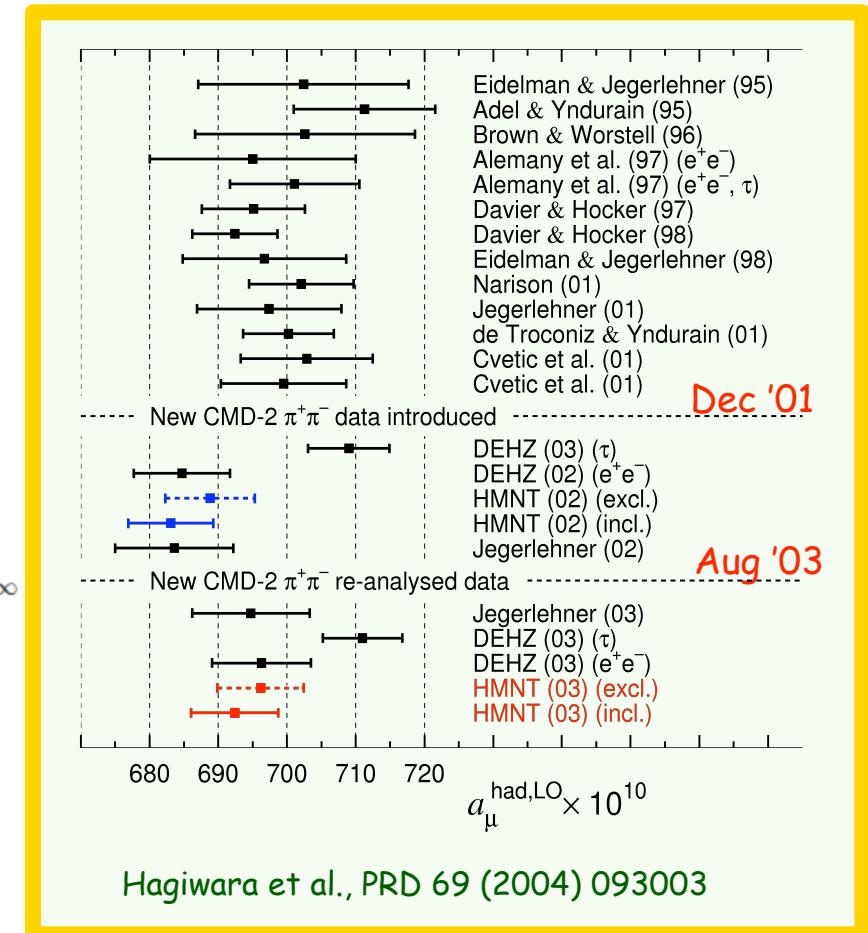
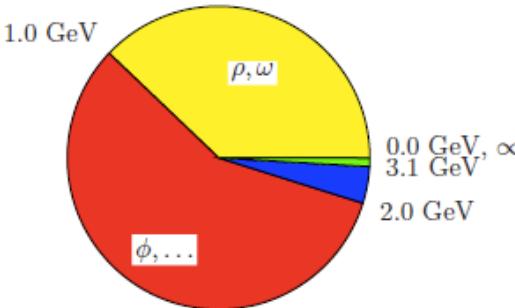
$$K(s) = \int_0^1 dx \frac{x^2(1-x)}{x^2 + (1-x)s/m_\mu^2}$$

Bouchiat & Michel 1961; Gourdin & de Rafael 1969

Central values



Errors²



F. Jegerlehner and A. Nyffeler, arXiv:0902.3360

The HLO contribution: e^+e^- data

$$a_\mu^{\text{HLO}} = 6909 \text{ (39)}_{\text{exp}} \text{ (19)}_{\text{rad}} \text{ (7)}_{\text{qcd}} \times 10^{-11} \quad \text{S. Eidelman, ICHEP06; M. Davier, TAU06}$$
$$= 6894 \text{ (42)}_{\text{exp}} \text{ (18)}_{\text{rad}} \times 10^{-11} \quad \text{Hagiwara, Martin, Nomura, Teubner, PLB649(2007)173}$$
$$= 6903 \text{ (53)}_{\text{tot}} \times 10^{-11} \quad \text{F. Jegerlehner, A. Nyffeler, arXiv:0902.3360}$$
$$= 6898 \text{ (43)}_{\text{exp+rad}} \text{ (7)}_{\text{qcd}} \times 10^{-11} \quad \text{Davier et al, arXiv:0906.5443v2 (all data)}$$

- 💡 **Radiative Corrections** (Luminosity, ISR, Vacuum Polarization, FSR) are a very delicate issue! Are they all under control?
- 💡 **CMD2:** The 1998 $\pi^+\pi^-$ data in the ρ energy range, published in 2007, agree with their earlier 1995 ones.
- 💡 **SND's** $\pi^+\pi^-$ 2006 data reanalysis in agreement with CMD2.

The HLO contribution: e^+e^- data (ISR Method)

- The RADIATIVE RETURN (ISR) Method: KLOE & BaBar.
Collider operates at fixed energy but s_π can vary continuously.
Important independent method made possible by beautiful
interplay between theory and experiment.
- KLOE: In 2008 KLOE presented an update of their 2005 $\pi^+\pi^-$
analysis plus a new measurement: PLB670 (2009) 285. The new
measurement supersedes the 2005 one.
- Agreement between KLOE (2008) and CMD2-SND below the ρ ,
some discrepancies above. Their contributions to a_μ^{HLO} agree.
- BaBar: $\pi^+\pi^-$ preliminary results (from 0.5 to 3 GeV) presented at
Tau08. Disagreement with CMD2, SND and KLOE. Better
agreement with τ results, especially with Belle. We must wait
for a publication.

The HLO contribution: Tau-decay data (Aleph, Opal, Cleo & Belle)

- The τ data of ALEPH and CLEO are significantly higher than the CMD2-SND-KLOE ones, particularly above the ρ .
- The 2008 $a_\mu^{\pi\pi} \tau$ result of Belle agrees with Aleph-Cleo-Opal. Some deviations from Aleph's spectral functions.
- Yesterday: $a_\mu^{\text{HLO}} = 7053 (45) \times 10^{-11}$ Davier et al, arXiv:[0906.5443v2](https://arxiv.org/abs/0906.5443v2)
- vs 2003: $a_\mu^{\text{HLO}} = 7110 (58) \times 10^{-11}$ Davier et al, EPJC31 (2003) 503

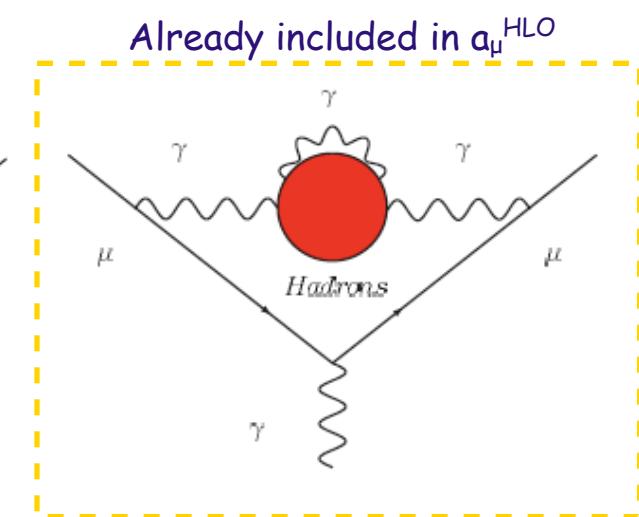
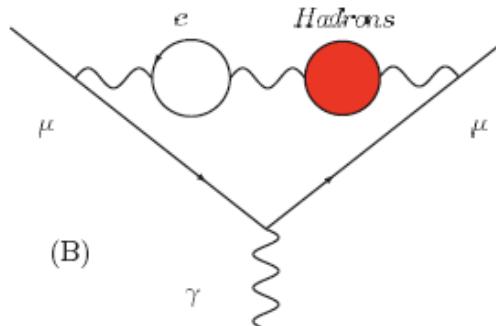
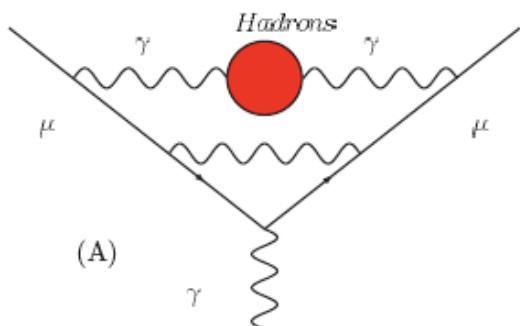
Belle's data included + IB corrections revisited & updated

(Marciano & Sirlin '88; Cirigliano, Ecker, Neufeld '01-'02, Flores-Baez et al. '06 & '07)

- The discrepancy with e^+e^- data is smaller, but it's still there! Inconsistencies in e^+e^- or τ data? All possible isospin-breaking (IB) effects taken into account? Recent claims that e^+e^- & τ data are consistent after IB effects & vector meson mixings considered (Benayoun et al.'07 & '09).

The hadronic higher-order (HHO) contributions: VP

• HHO: Vacuum Polarization



$O(\alpha^3)$ contributions of diagrams containing hadronic vacuum polarization insertions:

$$a_\mu^{\text{HHO(vp)}} = -98 (1) \times 10^{-11}$$

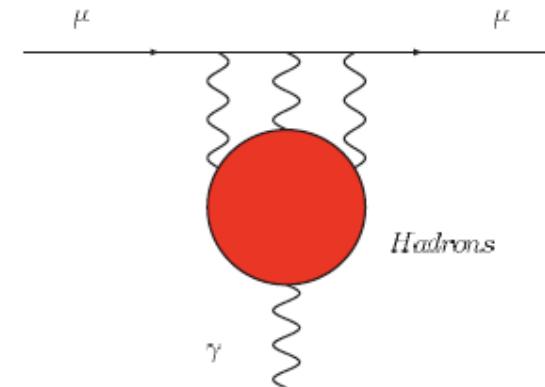
Krause '96, Alemany et al. '98, Hagiwara et al. '03 & '06

Shifts by $\sim -3 \times 10^{-11}$ if τ data are used instead of the e^+e^- ones
Davier & Marciano '04.

The hadronic higher-order (HHO) contributions: LBL

HHO: Light-by-light contribution

- Unlike the HLO term, for the hadronic l-b-l term one must rely on theoretical approaches.
- This term had a **troubled life!** Its recent determinations vary between:



$$a_\mu^{\text{HHO}}(|\mathbf{b}|) = +80 \text{ (40)} \times 10^{-11} \quad \text{Knecht \& Nyffeler '02}$$

$$a_\mu^{\text{HHO}}(|\mathbf{b}|) = +136 \text{ (25)} \times 10^{-11} \quad \text{Melnikov \& Vainshtein '03}$$

$$a_\mu^{\text{HHO}}(|\mathbf{b}|) = +105 \text{ (26)} \times 10^{-11} \quad \text{Prades, de Rafael, Vainshtein '09}$$

$$a_\mu^{\text{HHO}}(|\mathbf{b}|) = +116 \text{ (39)} \times 10^{-11} \quad \text{Jegerlehner \& Nyffeler '09}$$

(results based also on Hayakawa, Kinoshita '98 & '02; Bijnens, Pallante, Prades '96 & '02)

- Upper bound: $a_\mu^{\text{HHO}}(|\mathbf{b}|) < \sim 160 \times 10^{-11}$ Erler & Sanchez 2006, Pivovarov 2002
- Lattice? Very hard, but in progress: Rakow et al, Hayakawa et al., ...
- It's likely to become the ultimate limitation of the SM prediction.

The muon $g-2$: Standard Model vs. Experiment

- Adding up all the above contribution we get the following SM predictions for a_μ and comparisons with the measured value:

| $a_\mu^{\text{SM}} \times 10^{11}$ | $\Delta a_\mu \times 10^{11}$ | σ |
|------------------------------------|-------------------------------|----------|
| [1] 116 591 788 (51) | 292 (81) | 3.6 |
| [2] 116 591 773 (53) | 307 (82) | 3.7 |
| [3] 116 591 782 (59) | 298 (86) | 3.4 |
| [4] 116 591 777 (51) | 303 (81) | 3.7 |
| [5] 116 591 929 (52) | 151 (82) | 1.8 |

with $a_\mu^{\text{HHO}}(|b|) = 105 (26) \times 10^{-11}$. $\Delta a_\mu = a_\mu^{\text{EXP}} - a_\mu^{\text{SM}}$.

- [1] Eidelman at ICHEP06 & Davier at TAU06.
- [2] Hagiwara, Martin, Nomura, Teubner, PLB649 (2007) 173.
- [3] F. Jegerlehner and A. Nyffeler, arXiv:0902.3360.
- [4] Davier et al, arXiv:0906.5443v2 August 2009 (all e^+e^- data).
- [5] Davier et al, arXiv:0906.5443v2 August 2009 (τ data).

- The th error is now about the same as the exp. one!

The muon $g-2$ and the bounds on the Higgs mass

MP, W.J. Marciano & A. Sirlin

[PRD78, 013009 (2008)]

The Hadronic Contribution to $\alpha(M_Z^2)$...

The effective fine-structure constant at the scale M_Z^2 is given by:

$$\alpha(M_Z) = \frac{\alpha}{1 - \Delta\alpha(M_Z)} \quad \text{with} \quad \Delta\alpha = \Delta\alpha_{\text{lep}} + \Delta\alpha_{\text{had}}^{(5)} + \Delta\alpha_{\text{top}}$$

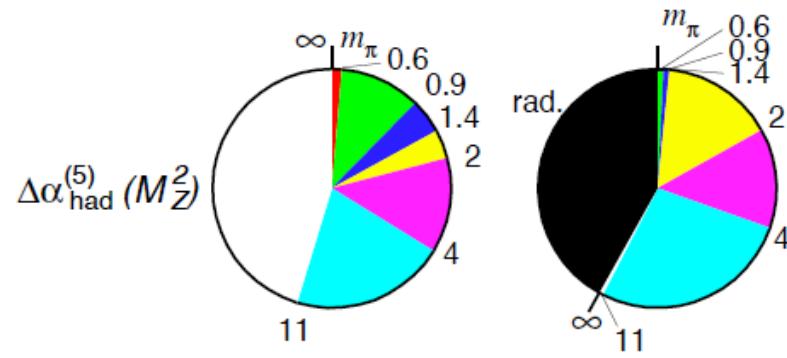
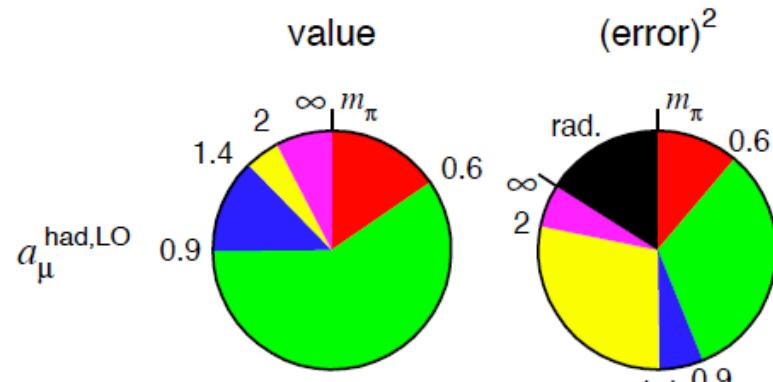
The light quarks part is determined by:

$$\Delta\alpha_{\text{had}}^{(5)}(M_Z) = \frac{M_Z^2}{4\alpha\pi^2} P \int_{4m_\pi^2}^\infty ds \frac{\sigma(s)}{M_Z^2 - s}$$

Progress due to significant improvement
of the data:

$$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) =$$

| | |
|--------------|--------------------------|
| 0.02800 (70) | Eidelman, Jegerlehner'95 |
| 0.02775 (17) | Kuhn, Steinhauser 1998 |
| 0.02749 (12) | Troconiz, Yndurain 2005 |
| 0.02758 (35) | Burkhardt, Pietrzyk 2005 |
| 0.02768 (22) | Hagiwara et al. 2006 |
| 0.02761 (23) | F. Jegerlehner 2008 |



Hagiwara et al (HMNT) 2006

... and the EW Bounds on the SM Higgs mass

- The dependence of SM predictions on the Higgs mass, via loops, provides a powerful tool to set bounds on its value.
- Comparing the theoretical predictions of M_W and $\sin^2\theta_{\text{eff}}^{\text{lept}}$
[convenient formulae in terms of M_H , M_{top} , $\Delta\alpha_{\text{had}}^{(5)}(M_Z)$ and $\alpha_s(M_Z)$ by Degrassi, Gambino, MP, Sirlin '98; Degrassi, Gambino '00; Ferroglio, Ossola, MP, Sirlin '02; Awramik, Czakon, Freitas, Weiglein '04 & '06]

with

$$M_W = 80.399 \text{ (25) GeV} \quad [\text{LEP+Tevatron}]$$
$$\sin^2\theta_{\text{eff}}^{\text{lept}} = 0.23153 \text{ (16)} \quad [\text{LEP+SLC}]$$

and

$$\Delta\alpha_{\text{had}}^{(5)}(M_Z) = 0.02768 \text{ (22)} \quad [\text{HMNT '07}]$$
$$M_{\text{top}} = 172.4 \text{ (1.2) GeV} \quad [\text{CDF-D0, Aug '08}]$$
$$\alpha_s(M_Z) = 0.118 \text{ (2)} \quad [\text{PDG '08}]$$

we get

$$M_H = 88^{+32}_{-24} \text{ GeV} \quad \& \quad M_H < 145 \text{ GeV } 95\% \text{ CL}$$

- The value of $\Delta\alpha_{\text{had}}^{(5)}(M_Z)$ is a key input of these EW fits...

Back to Δa_μ : how do we explain it?

- Δa_μ can be explained in many ways: errors in HHO-LBL, QED, EW, HHO-VP, g-2 EXP, HLO; or New Physics.
- Can Δa_μ be due to hypothetical mistakes in the hadronic $\sigma(s)$?
- An upward shift of $\sigma(s)$ also induces an increase of $\Delta \alpha_{\text{had}}^{(5)}(M_Z)$.
- Consider:

$$a_\mu^{\text{HLO}}: \quad a = \int_{4m_\pi^2}^{s_u} ds f(s) \sigma(s), \quad f(s) = \frac{K(s)}{4\pi^3}, \quad s_u < M_Z^2,$$

$$\Delta \alpha_{\text{had}}^{(5)}: \quad b = \int_{4m_\pi^2}^{s_u} ds g(s) \sigma(s), \quad g(s) = \frac{M_Z^2}{(M_Z^2 - s)(4\alpha\pi^2)},$$

and the increase

$$\Delta \sigma(s) = \epsilon \sigma(s)$$

($\epsilon > 0$), in the range:

$$\sqrt{s} \in [\sqrt{s_0} - \delta/2, \sqrt{s_0} + \delta/2]$$

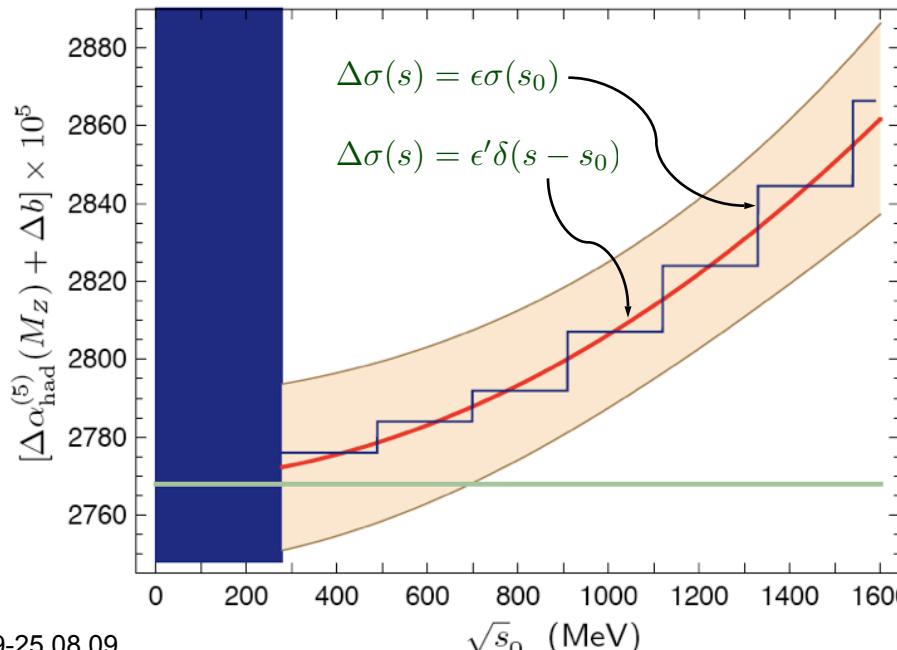


Shifts of a_μ^{HLO} and $\Delta\alpha_{\text{had}}^{(5)}(M_Z)$

- If this shift $\Delta\sigma(s)$ in $[\sqrt{s}_0 - \delta/2, \sqrt{s}_0 + \delta/2]$ is adjusted to bridge the $g-2$ discrepancy, the value of $\Delta\alpha_{\text{had}}^{(5)}(M_Z)$ increases by:

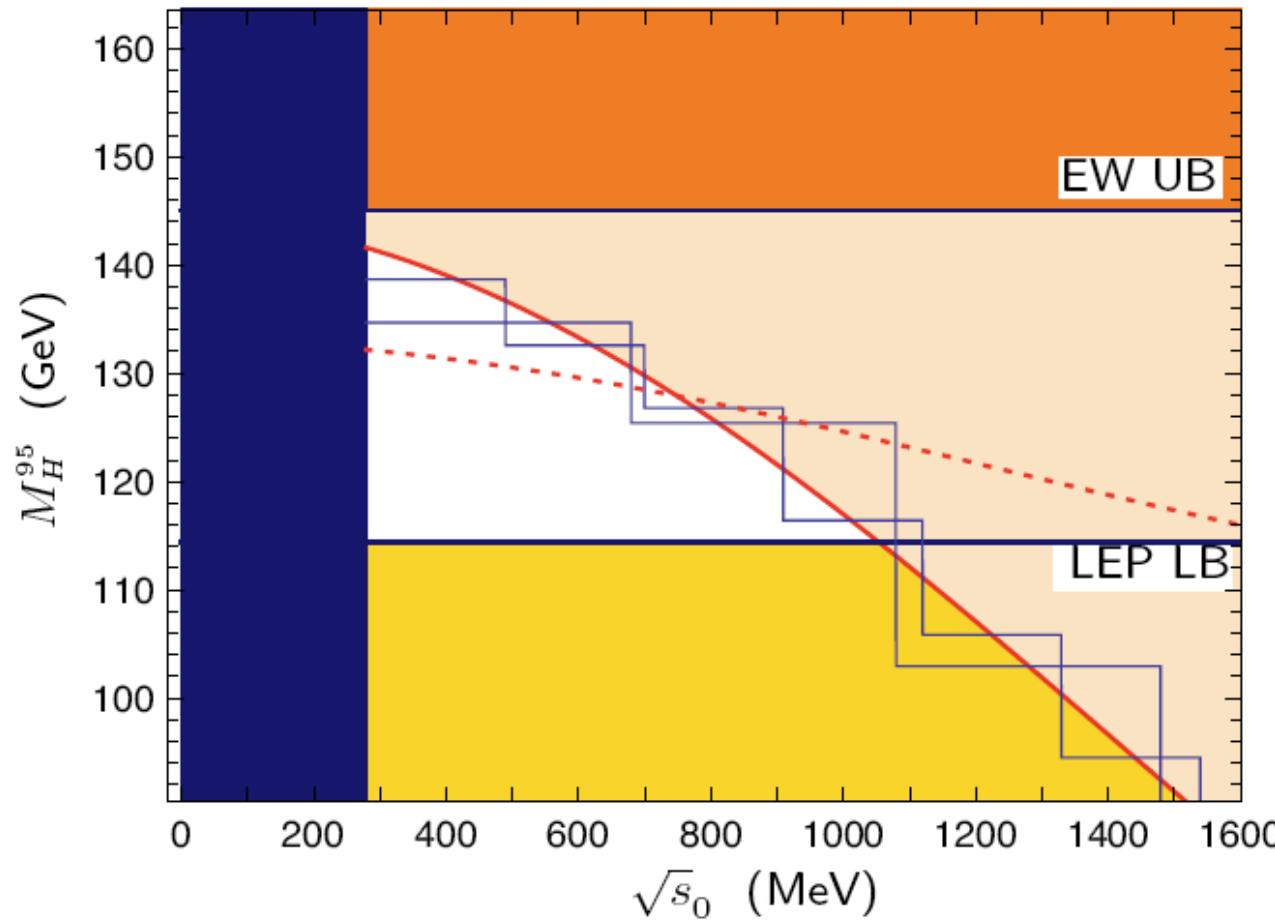
$$\Delta b(\sqrt{s}_0, \delta) = \Delta a_\mu \frac{\int_{\sqrt{s}_0 - \delta/2}^{\sqrt{s}_0 + \delta/2} g(t^2) \sigma(t^2) t dt}{\int_{\sqrt{s}_0 - \delta/2}^{\sqrt{s}_0 + \delta/2} f(t^2) \sigma(t^2) t dt}$$

- Adding this shift to $\Delta\alpha_{\text{had}}^{(5)}(M_Z) = 0.02768(22)$ [HMNT07], with $\Delta a_\mu = 302(88) \times 10^{-11}$ [HMNT07], we obtain:



The muon $g-2$: connection with the SM Higgs mass

- How much does the M_H upper bound change when we shift $\sigma(s)$ by $\Delta\sigma(s)$ [and thus $\Delta\alpha_{\text{had}}^{(5)}(M_Z)$ by Δb] to accommodate $\Delta\alpha_\mu$?

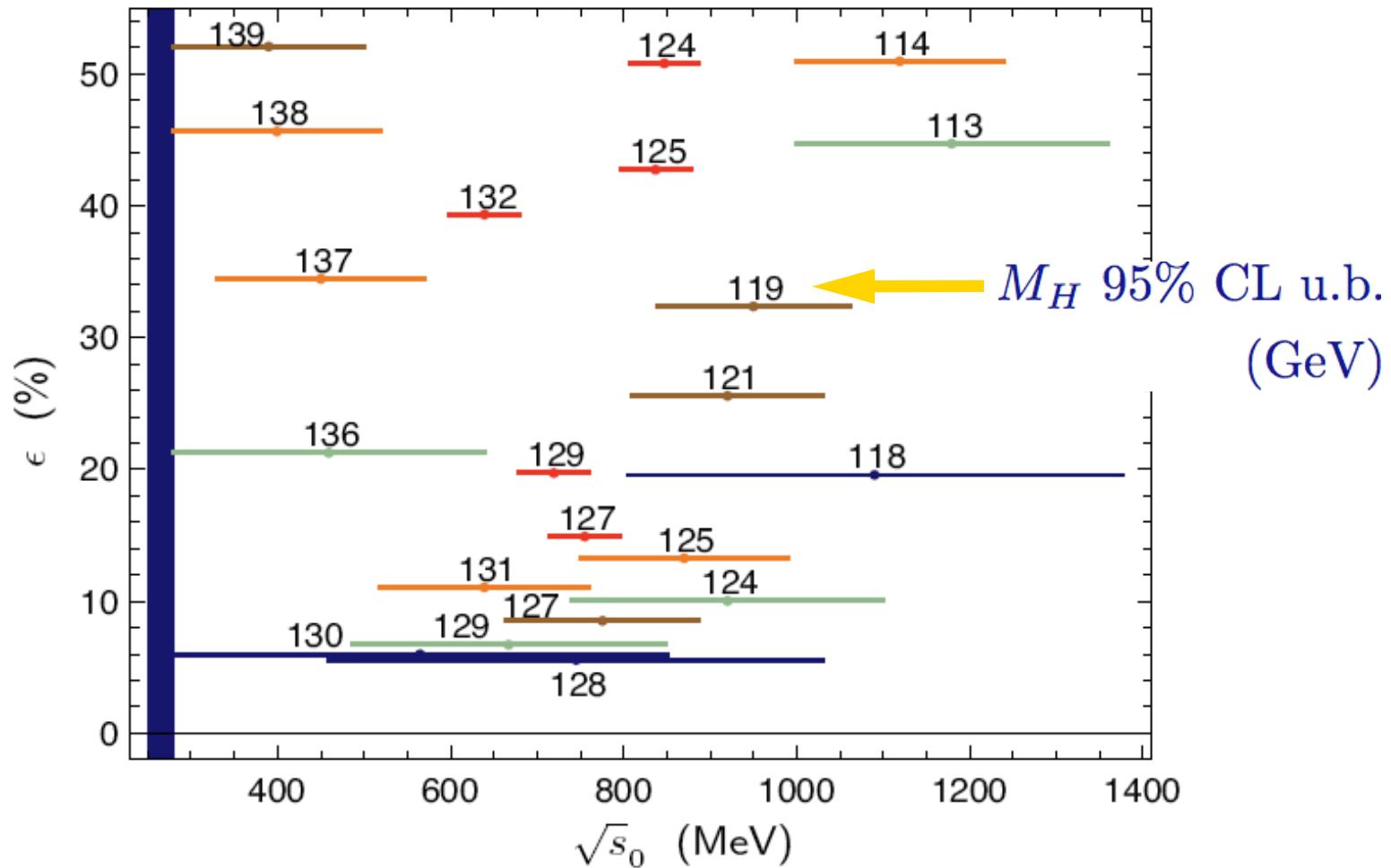


The muon g-2: connection with the SM Higgs mass (2)

- The LEP direct-search lower bound is $M_H^{LB} = 114.4 \text{ GeV}$ (95%CL).
- The hypothetical shifts $\Delta\sigma = \varepsilon\sigma(s)$ that bridge the muon g-2 discrepancy conflict with the LEP lower limit when $\sqrt{s_0} > \sim 1.2 \text{ GeV}$ (for bin widths δ up to several hundreds of MeV).
- While the use of tau data in the calculation of a_μ^{HLO} reduces the muon g-2 discrepancy, it increases the value of $\Delta a_{\text{had}}^{(5)}(M_Z)$, lowering the M_H upper bound (tension with the M_H lower bound).
- Recent claim: e^+e^- & tau data consistent below $\sim 1 \text{ GeV}$ (after isospin viol. effects & vector meson mixings). We could thus assume that Δa_μ is fixed by hypothetical errors above $\sim 1 \text{ GeV}$ (where disagreement persists). If so, M_H^{UB} falls below M_H^{LB} !!
- Scenarios where Δa_μ is accommodated without affecting M_H^{UB} are possible, but considerably more unlikely.

How realistic are these shifts $\Delta\sigma(s)$?

- How realistic are these shifts $\Delta\sigma(s)$ when compared with the quoted exp. uncertainties? Study the ratio $\varepsilon = \Delta\sigma(s)/\sigma(s)$:



How realistic are these shifts $\Delta\sigma(s)$? (2)

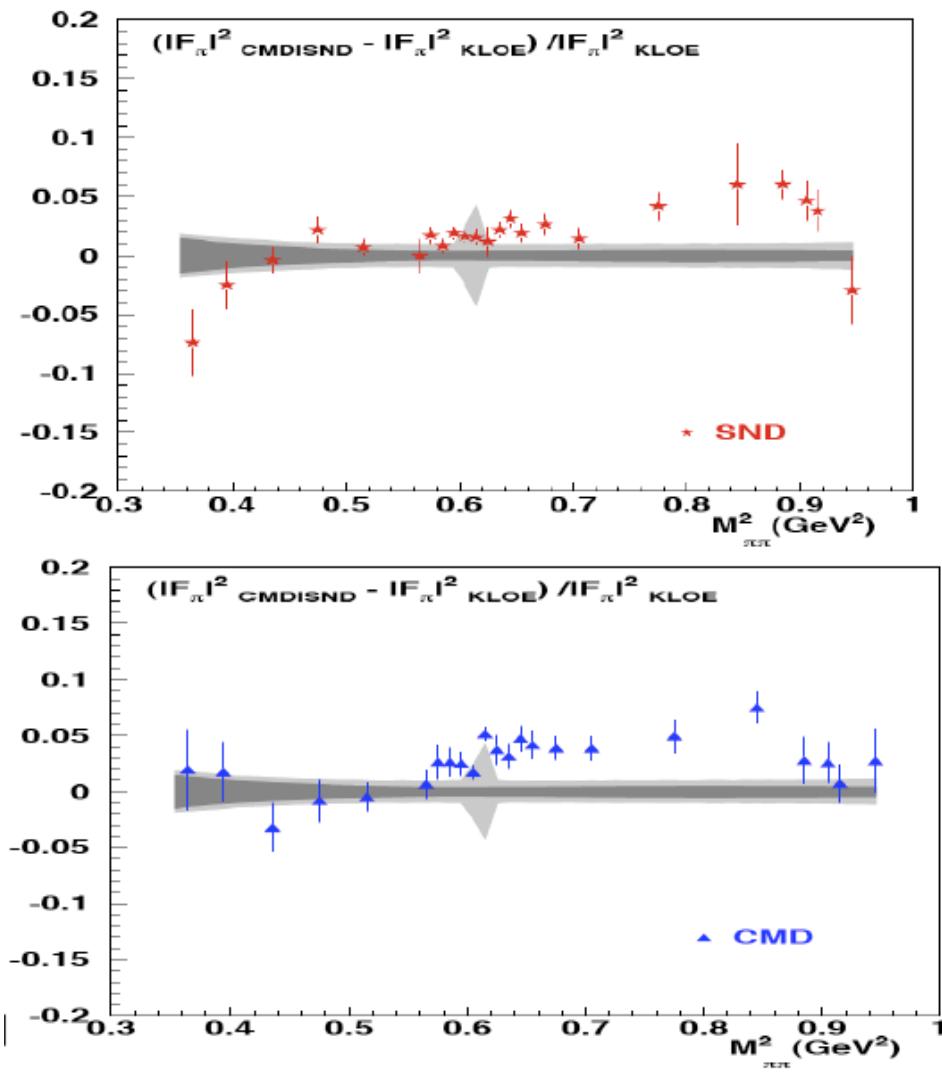
- The minimum ε is $\sim +4\%$. It occurs if σ is multiplied by $(1+\varepsilon)$ in the whole integration region (!), leading to $M_H^{UB} \sim 70 \text{ GeV}$ (!!)
- As the quoted exp. uncertainty of $\sigma(s)$ below 1 GeV is \sim a few per cent (or less), the possibility to explain the muon g-2 with these shifts $\Delta\sigma(s)$ appears to be unlikely.
- If, however, we allow variations of $\sigma(s)$ up to $\sim 6\%$ (7%), M_H^{UB} is reduced to less than $\sim 130 \text{ GeV}$ (131 GeV). E.g., the $\sim 6\%$ shift in the interval $[0.6, 1.2] \text{ GeV}$, required to fix Δa_μ , lowers M_H^{UB} to 126 GeV. Tension with the $M_H > \sim 120 \text{ GeV}$ "vacuum stability" bound.
- Reminder: the above M_H upper bounds, like the LEP-EWWG ones, depend on the value of $\sin^2 \theta_{\text{eff}}^{\text{lept}}$. They also depend on M_t & its unc. δM_t . We prepared simple formulae to translate easily M_H upper bounds discussed above into new values corresponding to M_t & δM_t inputs different from those employed here.

Conclusions

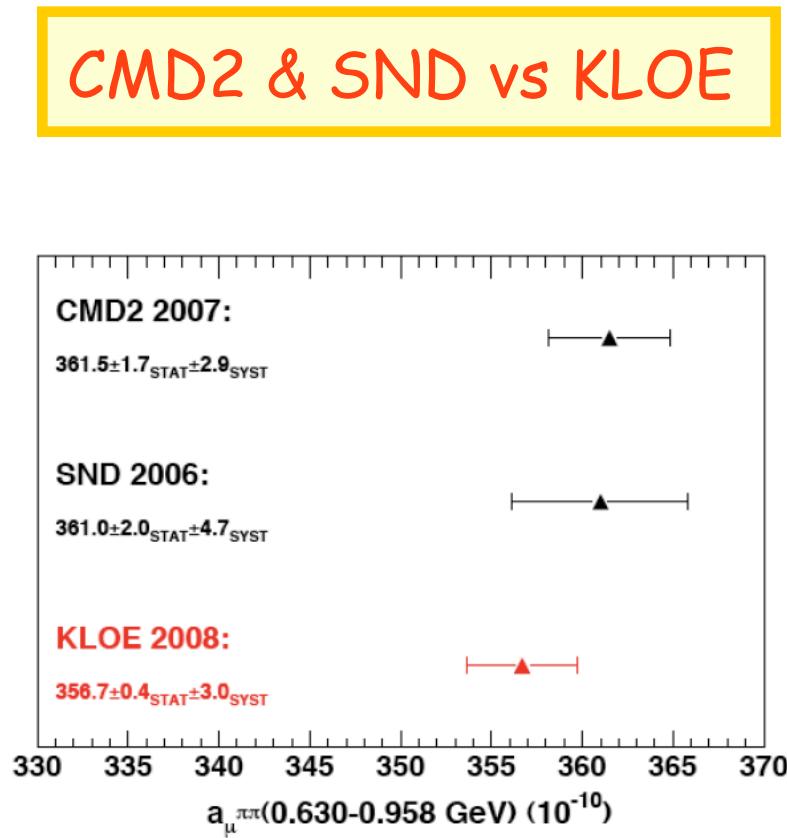
- g : Beautiful examples of interplay between theory and experiment:
 g_e probed at eV → α and extraordinary test of QED's validity;
 g_μ probed at ppb → test of the full SM and great opportunity to unveil (or just constrain) "New Physics" effects!
- The discrepancy Δa_μ is more than 3σ if e^+e^- data are used. With tau data, the deviation is $\sim 2\sigma$. BaBar 2π ? More e^+e^- data & analyses eagerly awaited! QED & EW ready for new $g-2$ exp! LBL??
- Δa_μ can be due to New Physics, or to problems in a_μ^{SM} (or a_μ^{EXP}). Can it be due to hypothetical mistakes in the hadronic $\sigma(s)$? An increase $\Delta\sigma(s)$ could bridge Δa_μ , leading however to a decrease on the EW upper bound on the SM Higgs mass M_H ...
- By means of a detailed analysis we conclude that solving Δa_μ via an increase of $\sigma(s)$ is unlikely in view of current exp. error estimates. However, if this turns out to be the solution, then the M_H upper bound drops to about 130 GeV which, in conjunction with the LEP 114 GeV direct lower limit, leaves a rather narrow window for M_H .

The End

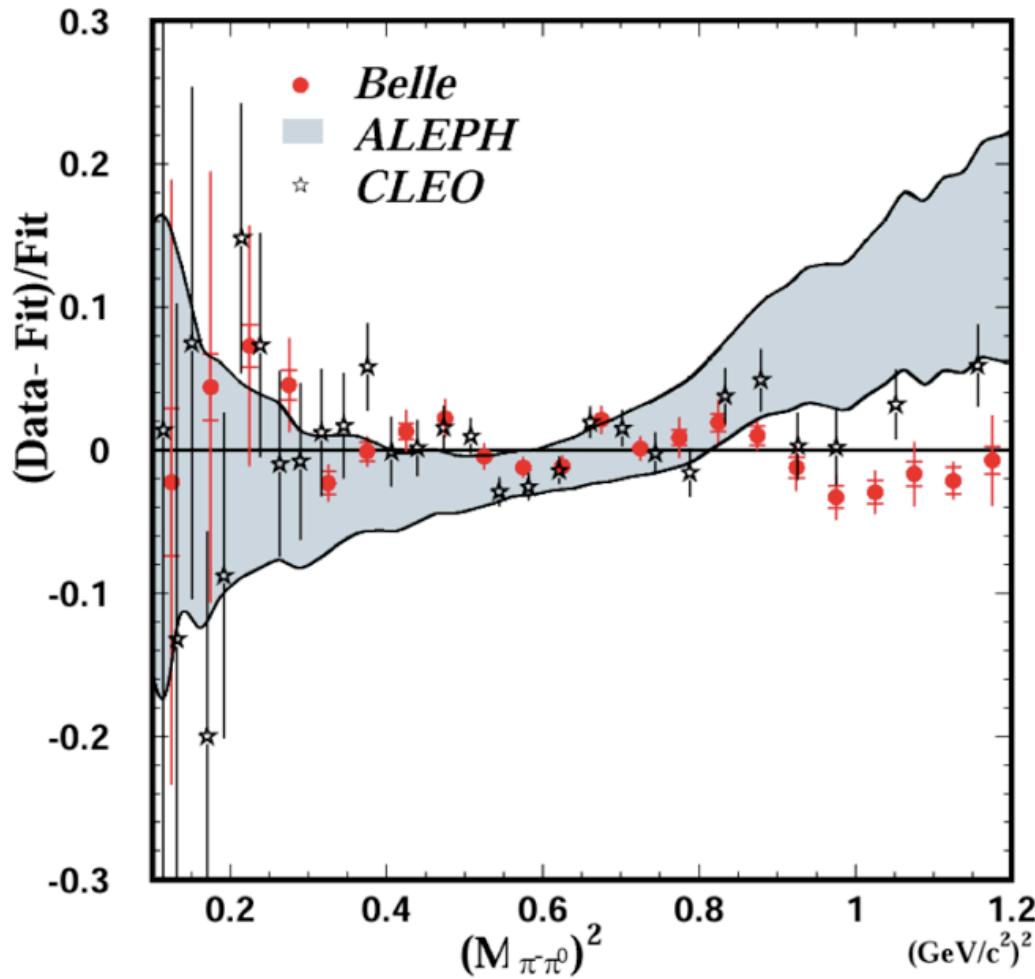
Back-up Slides



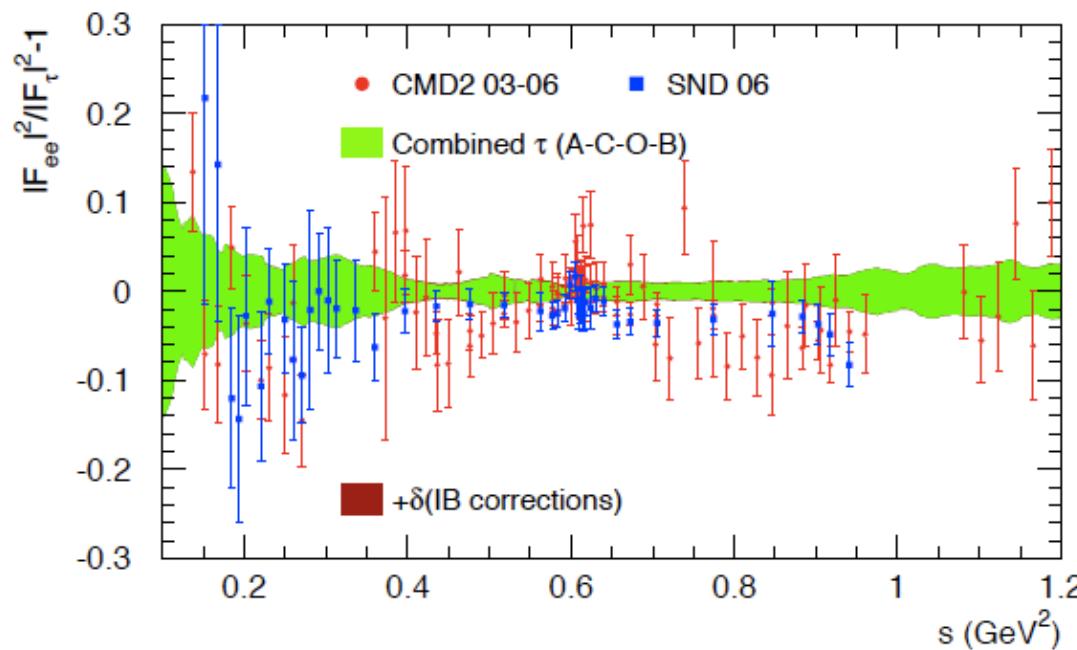
band: KLOE error
data points: CMD2/SND experiments



G. Venanzoni, Tau08, Novosibirsk,
September 2008



Fujikawa, Hayashii, Eidelman [for the Belle Collab.], arXiv:0805.3773, May '08



Davier et al,
arXiv:0906.5443v2

